

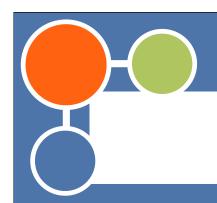
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Work in collaboration with Tilmino sPlehn

FNAL Workshop September 2008



### Outline



- What the heck is a sgluon?
  - O The MRSSM
- Sgluon couplings
- Sgluon production and decay
- Signals at the LHC
- Outlook





## What is a sgluon?



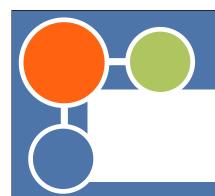
- A sgluon is a scalar gluon.
  - An electrically neutral, color octet, spin 0 particle.
- Such objects have been studied in the past as arising from two (or more) UEDs - "spinless adjoints".

Dobrescu, Kong, Mahbubani

I'm interested in a different context, with different phenomenology. Sgluons are extra states that appear in a supersymmetric theory with a Dirac gluino mass.

"Supersoft" - Fox, Nelson, Weiner





### **MRSSM**



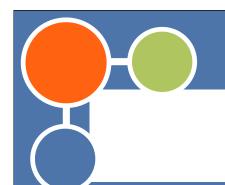
- Last year, an R-symmetric version of the minimal supersymmetric Standard Model appeared.
  Kribs, Poppitz, Weiner
- The R-symmetry turns out to be very useful in order to help with the SUSY flavor problem.

$$\frac{1}{m_{\tilde{g}}} \tilde{s}_L^* \tilde{d}_R \quad \bar{d}_R s_L \quad \longrightarrow \frac{1}{m_{\tilde{g}}^2} \tilde{s}_L^* \partial_\mu \tilde{d}_L \quad \bar{d}_L \gamma^\mu s_L$$
 Dimension 5 Dimension 6

- Relatively flavor-mixed squark masses become compatible with measurements of flavor-violation.
- Also see Andrew's talk coming up later this week!



Amigo, Blechman, Fox, Poppitz



## Dirac Gauginos



- Majorana gaugino masses violate the R-symmetry.
- The solution is to enlarge the theory to include enough. degrees of freedom for a Dirac gaugino by adding adjoint chiral super-multiplets.  $R_{\Phi}=0$

$$\Phi^a = \phi^a + \sqrt{2}\theta\psi^a + \theta^2 F_{\Phi}$$

Fox, Nelson, Weiner

A D-term source of SUSY-breaking results in Dirac  $\int d^2\theta \frac{\sqrt{2}}{M} W' W_3^a \Phi^a \to m_{\tilde{g}} \lambda^a \psi^a + \dots$  $\langle W' \rangle = D\theta \qquad m_{\tilde{g}} = \frac{D}{M}$ gaugino masses:

$$\langle W' \rangle = D\theta \qquad m_{\tilde{g}} = \frac{D}{M}$$



## SUSY Breaking



- Since Φ contains a complex scalar, technically we actually have two (real) sgluons.
- They gain mass directly from either D-term or F-term SUSY-breaking spurions:  $\langle X \rangle = \theta^2 F$

$$\int d^4\theta \left\{ \frac{1}{M_1^2} X^{\dagger} X \Phi^{\dagger} \Phi + \frac{1}{M_2^2} X^{\dagger} X \Phi^2 \right\} + \int d^2\theta \frac{1}{M_3^2} W_{\alpha}' W'^{\alpha} \Phi^2 + H.c.$$

These operators induce both  $\varphi^2$  as well as  $|\varphi|^2$  mass terms, generically splitting the complex scalar into two real mass eigenstates.



#### **D**-terms



The gluino mass operator also contributes through the SU(3) D-term (remember the ...s):

$$\int d^2\theta \frac{\sqrt{2}}{M} W' W_3^a \Phi^a \to \dots + \sqrt{2} D^a \left( m_{\tilde{g}} \phi^a + m_{\tilde{g}}^* \phi^{a*} \right)$$

Going on-shell (integrating out the D-terms) produces more of the same types of masses, and interactions:

$$\frac{1}{2}D^{a}D^{a} + \sqrt{2}D^{a}\left(m_{\tilde{g}}\phi^{a} + m_{\tilde{g}}^{*}\phi^{a*}\right) + D^{a}g_{S}\sum\tilde{q}^{*}T^{a}\tilde{q}$$

$$\rightarrow m_{\tilde{g}}^{2}\phi^{2} + m_{\tilde{g}}^{2*}\phi^{*2} + |m_{\tilde{g}}|^{2}|\phi|^{2} + \sqrt{2}g_{S}\left(m_{\tilde{g}}\phi^{a} + m_{\tilde{g}}^{*}\phi^{a*}\right)\left(\sum\tilde{q}^{*}T^{a}\tilde{q}\right) + \dots$$
Usual SU(3) D-terms

More sgluon mass terms

Interactions with squarks

## Sgluon Masses



- We've learned that sgluons typically get masses from SUSY-breaking of order the gluino mass.
- However, I can play the other soft masses against those contributions and really end up with anything I want. So I will lump them all together and call them parameters in their own right:

$$\mathcal{L}_{mass} = -m_1^2 |\phi^a|^2 - \frac{1}{2} |m_2^2| e^{i\gamma} \phi^{a2} - \frac{1}{2} |m_2^{*2}| e^{-i\gamma} \phi^{a*2}$$

The mass eigenvalues are:  $m_{\phi_1,\phi_2}^2 = m_1^2 \mp |m_2^2|$ 

$$\phi_1^a = \sin\frac{\gamma}{2} \phi^a + \cos\frac{\gamma}{2} \phi^{a*}$$

l'll just take  $\gamma = 0$  from here on...



$$\phi_2^a = \cos \frac{\gamma}{2} \phi^a - \sin \frac{\gamma}{2} \phi^{a*}$$

# Sgluon Couplings

 Sgluons are color octets, and thus have interactions with gluons entirely determined by SU(3)<sub>C</sub> (SUSY) gauge invariance:

$$(D_{\mu}\phi)^* (D^{\mu}\phi) + i\sqrt{2}g_S f_{bc}^a \tilde{g}^b (\phi^a P_L + \phi^{a*} P_R) \tilde{g}^c$$

- O Thus, they can be pair-produced in a model-independent way at the LHC.
- No renormalizable supersymmetric interactions can couple the sgluon to matter fields.

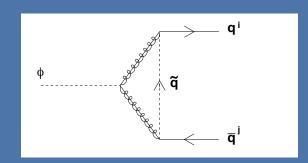


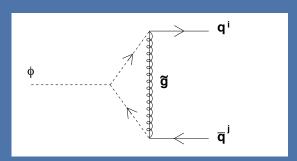
# Coupling to Quarks

- Being R-charge 0, nothing stops the sgluons from coupling to pairs of ordinary quarks.
- However, the gauge-invariant interaction is a higher dimensional operator:

$$\phi^a \left[ (H\bar{q}_L) T^a q_R \right] + H.c.$$

- Loops containing gluinos and squarks will induce this operator in the MRSSM.
- Chirality demands a gluino mass insertion in the first graph. The second graph is already proportional to the gluino mass, as we already saw.







# Coupling to Quarks

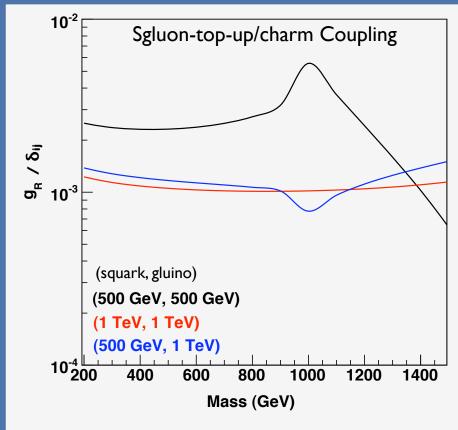


- We need one quark mass insertion. Interactions involving at least one top are preferred.
- $\sqrt{2}g_S^3 N_C m_{\tilde{g}} m_t C_{11} \left( m_t^2, 0, M^2; m_{\tilde{g}}, m_{\tilde{q}}, m_{\tilde{g}} \right)$

$$-\sqrt{2}g_S^3 \frac{1}{N_C} m_{\tilde{g}} m_t \left[ C_{11} \left( m_t^2, 0, M^2; m_{\tilde{q}}, m_{\tilde{g}}, m_{\tilde{q}} \right) + C_0(...) \right]$$

- The other quark can be either up or charm, with no penalty other than the squark mixing parameters.
- In the MRSSM we expect reasonably large squark mixing, so relatively large coupling to:







## Sgluon Decays



- Thus, sgluons are going to prefer to decay into top quarks.
- Assuming (in the spirit of the MRSSM) that the upcharm-top mixings are all large, there should be comparable branching ratios into top + up or charm as well.
  Flavor-conserving pheno: Dobrescu, Kong, Mahbubani
- Sgluons carry no charge or fermion number. They will have large branching ratios into:

$$t\bar{t}, t\bar{u}, t\bar{c}, u\bar{t}, c\bar{t}$$

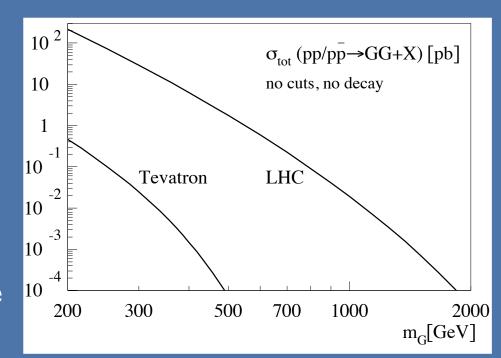




#### Pair Production



- Pair production rates depend only on the mass of the sgluon (the couplings are gs).
- There are events for masses up to about 400 GeV (Tevatron) and 2 TeV (LHC).
- Single production from quarks is extremely tiny because of the small, loop-induced coupling strength.



A promising signature: 2 like-sign tops (with 2 jets) with the tops decaying to e or  $\mu$ .



(Arvind told us about this class of signatures yesterday!)



## **Tevatron Limits**





Bar-Shalom, Rajaraman, Whiteson, Yu

When he suggested we take out our iPhones and photograph the CDF limits, I'm not sure he really thought anyone would take the invitation literally...





- We are still finalizing our analysis.
- Options include:
  - Flavor-violating decays, resulting in two tops and two jets.
    - Like-sign tops revealed through leptonic top decays
    - Lepton+jets top decays to reconstruct at least one sgluon.
  - O Flavor-conserving decays into four tops total.
    - Again, like-sign top signatures.

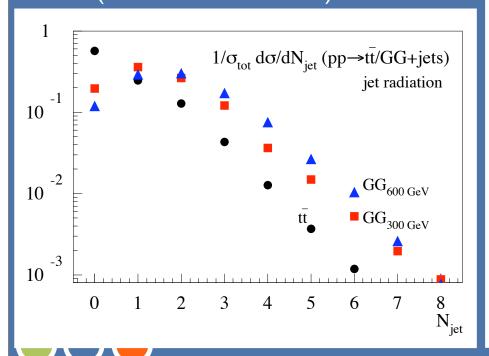
Also from top compositeness: Lillie, Shu, Tait

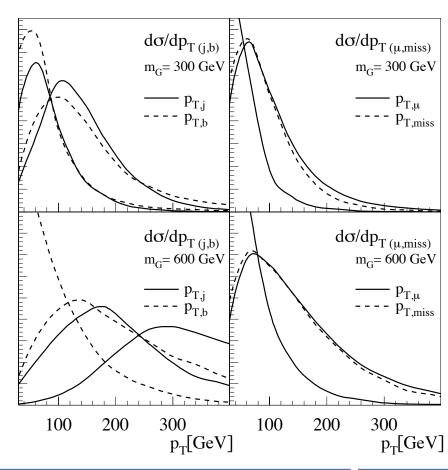
Sgluon masses up to order I TeV should be discoverable.

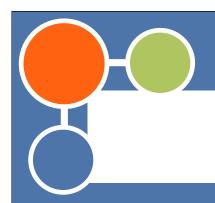


# Signal and Background

 We include the effects of QCD radiation through MLM matching to signal and t tbar background. (Thanks MadEvent...)







### Outlook



- Sgluons are cool!
- They are a generic signature of a model with Dirac gaugino masses, and an essential feature of the MRSSM.
- They can have highly flavor-violating couplings, and prefer to decay into at least one top.
  - O Their Branching Ratios tell us something about sfermion mixing.
- As color octets, there is large QCD pair production, resulting in like-sign top quarks at the LHC.

